

Precision Higgs Width at the ILC

Joshua Sayre

Based on Work with Tao Han and Zhen Liu

University of Pittsburgh

April 4, 2013

Table of contents

- 1 General Case for Γ_{Higgs} at the ILC
- 2 Strategies for Width Determination
- 3 Width at 250 GeV
- 4 Width at Design Energy

Where Do We Go From Here?

Prediction:

Junping will have already covered much of this.

We have a Higgs! (Or an unreasonable facsimile thereof...)

The LHC will continue to search for new physics in the form of SUSY, Extra Dimensions, Dark Matter, Additional Higgs, Z' , W' , 4^{th} Generation, etc., etc.

If nothing else it will continue to refine our knowledge of the Higgs via coupling strengths, mass and width.

So far (mostly) consistent with the Standard Model, but new physics may show up in precision measurements.

Where Do We Go From Here?

Prediction:

Junping will have already covered much of this.

We have a Higgs! (Or an unreasonable facsimile thereof...)

The LHC will continue to search for new physics in the form of SUSY, Extra Dimensions, Dark Matter, Additional Higgs, Z' , W' , 4^{th} Generation, etc., etc.

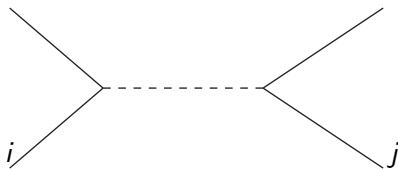
If nothing else it will continue to refine our knowledge of the Higgs via coupling strengths, mass and width.

So far (mostly) consistent with the Standard Model, but new physics may show up in precision measurements.

But there is a problem.

The Problem

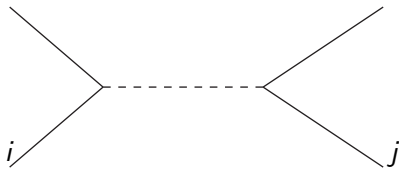
We don't really measure g_i and Γ at the LHC. We measure processes of the form



$$\sigma_{ij} \propto \frac{g_i^2 g_j^2}{\Gamma}$$

The Problem

We don't really measure g_i and Γ at the LHC. We measure processes of the form



$$\sigma_{ij} \propto \frac{g_i^2 g_j^2}{\Gamma}$$

$$\forall i : g_i \rightarrow x g_i \quad \Gamma \rightarrow x^4 \Gamma$$

Regardless of how many σ_{ij} s we measure.

Constraints are model dependent.

The Solution

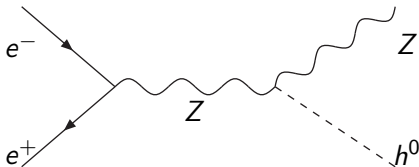
Measure $ii \rightarrow h \propto g_i^2$!

Break the degeneracy by measuring inclusive rates.

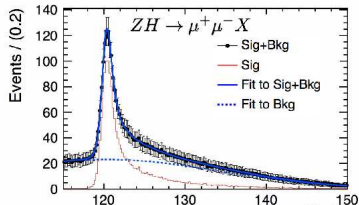
- But, virtually no way to do this at the LHC without assumptions.
- Can't rely on decay products.
- Instead use other particles in interaction to reconstruct a typical signature.
- Recoil mass to the rescue: $ii \rightarrow h + X$ then $M(ii - X) = M_h$
- Don't know initial states completely in a hadron collider, but we do at a lepton collider.

Recoil Mass at the ILC

At 250 GeV, we maximize Higgs production via “Higgstrahlung”.



This allows a sharp mass peak (especially with $Z \rightarrow \mu\mu$) from which we can determine the cross-section $\delta\sigma_{Zh} \sim 2.5\%$. [Hengne Li, 2009]



Higgs Width Determination in General

By measuring the inclusive cross-section $\sigma_{Zh} \equiv \sigma(e^+e^- \rightarrow Zh)$ we can extract the coupling constant g_{hZZ} (to $\sim 1\%$).

We can use this to get the total width, and with that we can normalize all other couplings correctly and model independently. **All model independent width determinations depend on this measurement.**

The most straightforward way to calculate the total width:

$$\Gamma = \frac{\sigma_{Zh}^2}{\sigma_{ZZ}}$$

where $\sigma_{ZZ} \equiv \sigma(e^+e^- \rightarrow Zh \rightarrow ZZZ)$.

Unfortunately this is a rather hard measurement due to small cross-section.

For 250fb^{-1} at 250 GeV, $\delta_{ZZ} \equiv \frac{\Delta\sigma_{ZZ}}{\sigma_{ZZ}} \sim 20\%$. I.e. $\delta_\Gamma \sim 20 + \%$.

Generalizing

g_Z is our only independently measure coupling, so we need something proportional to $\frac{g_Z^n}{\Gamma}$. In practice $n = 4$ is our only choice. But, rather than:

$$\Gamma = \frac{\sigma_{ZH}^2}{\sigma_{ZZ}}$$

We can use

Generalizing

g_Z is our only independently measure coupling, so we need something proportional to $\frac{g_Z^n}{\Gamma}$. In practice $n = 4$ is our only choice. But, rather than:

$$\Gamma = \frac{\sigma_{ZH}^2}{\sigma_{ZZ}}$$

We can use

$$\Gamma = \frac{\sigma_{ZH}^2 \sigma_{ij}}{\sigma_{Zi} \sigma_{Zj}}$$

Where σ_{ij} is now any measured cross-section with a Higgs coupling to i at one vertex and j at the other.

Generalizing

g_Z is our only independently measure coupling, so we need something proportional to $\frac{g_Z^n}{\Gamma}$. In practice $n = 4$ is our only choice. But, rather than:

$$\Gamma = \frac{\sigma_{ZH}^2}{\sigma_{ZZ}}$$

We can use

$$\Gamma = \frac{\sigma_{ZH}^2 \sigma_{ij}}{\sigma_{Zi} \sigma_{Zj}}$$

Where σ_{ij} is now any measured cross-section with a Higgs coupling to i at one vertex and j at the other.

Or even a step further

$$\Gamma = \frac{\sigma_{ZH}^2 \sigma_{ki} \sigma_{jl}}{\sigma_{Zi} \sigma_{Zj} \sigma_{kl}}$$

Combinations of Measurements

We can also benefit from averaging over equivalent combinations of measurements

$$\Gamma = \sigma_{ZH}^2 \left(\frac{1}{\sigma_{ZZ}} \oplus \sum_{i,j \neq Z} \frac{\sigma_{ij}}{\sigma_{Zi} \sigma_{Zj}} \right)$$
$$\Gamma = \sigma_{ZH}^2 \frac{\left(\frac{1}{\sigma_{ZZ}} \frac{1}{\delta_{ZZ}^2} + \sum_{i,j \neq Z} \frac{\sigma_{ij}}{\sigma_{Zi} \sigma_{Zj}} \frac{1}{(\delta_{iZ}^2 + \delta_{jZ}^2 + \delta_{ij}^2)} \right)}{\frac{1}{\delta_{ZZ}^2} + \sum_{i,j \neq Z} \frac{1}{\delta_{iZ}^2 + \delta_{jZ}^2 + \delta_{ij}^2}}$$

(But be careful when propagating errors.)

Also note that for some i, j , multiple processes can give us σ_{ij} .

$$\sigma_{WZ} = \sigma(Zh \rightarrow ZWW) \oplus \sigma(WW \rightarrow \nu\nu h \rightarrow \nu\nu ZZ)$$

In principle, we have many handles, although only some will be useful in terms of significance.

Expected Precisions

Energy	250 GeV	250 GeV	500 GeV	500 GeV
Process	$e^+e^- \rightarrow Zh$	$e^+e^- \rightarrow Z\nu\nu$	$e^+e^- \rightarrow e^+e^-h$	$e^+e^- \rightarrow Z\nu\nu$
δ_{Zb}	1.1%	10.5%	1.8%	0.6%
δ_{ZW}	6.4%		9.2%	
δ_{Zc}	7.4%		12%	6.2%
$\delta_{Z\tau}$	4.2%		5.4%	14%
$\delta_{Z\gamma}$	29 – 38%		29 – 38%	20 – 26%
δ_{Zg}	9.1%		14%	4.1%
δ_{Wb}				0.6%
δ_{WW}				2.6%
δ_{Wc}				6.2%
$\delta_{W\tau}$				14.0%
δ_{Wg}				4.1%

Numbers taken from DBD. Thanks to the work of many authors!

Optimum Channels

Clearly, b and W modes are dominant when accessible (for SM-like Higgs). But, at 250 GeV only some channels are available. We want to make maximum use of the data in hand.

- σ_{ZZ} measurable but with poor precision expected. $\delta\Gamma \sim 20\%$.
- Need $\sigma_{ij} : i, j \neq Z$ for other routes. Durig et al. report σ_{Wb} measurable at 10% level. This is the only visible W -fusion channel at this energy, but it makes possible

$$\Gamma = \frac{\sigma_{ZH}^2 \sigma_{Wb}}{\sigma_{Zb} \sigma_{ZW}}$$
$$\delta\Gamma \simeq 13.4\%.$$

- This exhausts pure ILC options. However, we can turn to LHC data as well.

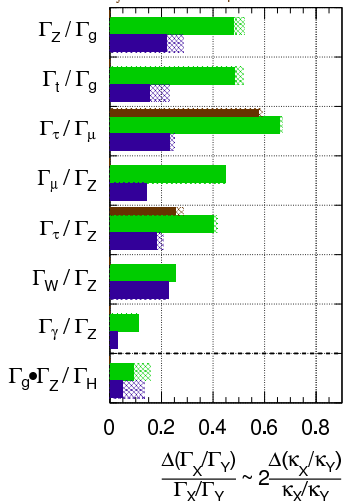
$$\Gamma = \frac{\sigma_{ZH}^2 \sigma_{Pi}}{\sigma_{PZ} \sigma_{Zi}}$$

LHC Branching Ratios, Atlas

ATLAS Preliminary (Simulation)

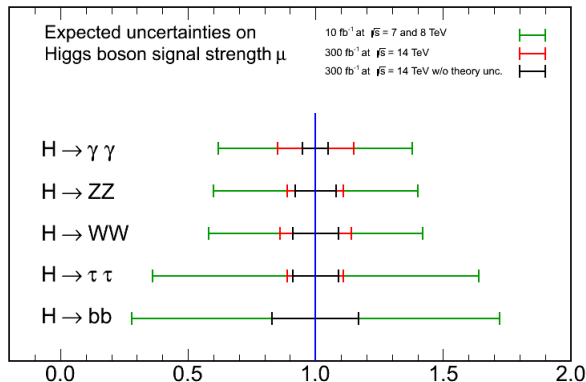
$\sqrt{s} = 14$ TeV: $\int \mathcal{L} dt = 300 \text{ fb}^{-1}$; $\int \mathcal{L} dt = 3000 \text{ fb}^{-1}$

$\int \mathcal{L} dt = 300 \text{ fb}^{-1}$ extrapolated from 7+8 TeV



LHC Higgs Couplings, CMS

CMS Projection



250 GeV cont.

Useful LHC channels:

LHC Meas.	LHC Error (300/3000)	$\delta\Gamma$
$\frac{\Gamma_\gamma}{\Gamma_Z}$	12%/3%	32%/30%
$\frac{\Gamma_W}{\Gamma_Z}$	26%/22%	27%/23.5%
$\frac{\Gamma_\tau}{\Gamma_Z}$	40(26)%/18%	40(27)%/19%
$\frac{\Gamma_b}{\Gamma_Z}$	20%/?	21%/ ?

Multiple constraints contribute to best Γ at this energy.

$$\frac{\Gamma}{\sigma_{Zh}^2} = \sigma_{ZZ} \oplus \frac{\sigma_{WZ}\sigma_{bZ}}{\sigma_{Wb}} \oplus \frac{\sigma_{WZ}\sigma_{PZ}}{\sigma_{PW}} \oplus \frac{\sigma_{\gamma Z}\sigma_{PZ}}{\sigma_{P\gamma}} \oplus \frac{\sigma_{\tau Z}\sigma_{PZ}}{\sigma_{P\tau}} \oplus \frac{\sigma_{bZ}\sigma_{PZ}}{\sigma_{Pb}}$$

But pay attention to polarizations!

Typical polarization in studies is $P_{e^-}P_{e^+} = (-0.8, 3.0)$ needed for WW fusion. But opposite polarization used in $Zh \rightarrow ZZZ, ZWW$ to suppress WW background.

Combined Result Estimates

Assuming numbers as given a combined analysis gives

$$\delta_{\Gamma} \simeq 10/9\%$$

However, the combination using $WW \rightarrow Zh \rightarrow Zbb$ uses left-handed polarized electrons. While the σ_{ZZ} and σ_{WZ} measurement rely on right-handed polarization to suppress background.

Using only LHC + right polarization

$$\delta_{\Gamma} \simeq 13/11\%$$

Options at 500 GeV

With full information from 500 GeV runs, the single best combination is likely to be

$$\Gamma = \sigma_{Zh}^2 \frac{\sigma_{Wb}^2}{\sigma_{Zb}^2 \sigma_{WW}}$$

- With 500fb^{-1} at 500 GeV, expected relative error $\delta\Gamma \simeq 11\%$.
- Assuming full information from 250 and 500 GeV runs $\delta\Gamma \simeq 6\%$.
- However, the more direct route

$$\Gamma = \frac{\sigma_{ZH}^2 \sigma_{Wb}}{\sigma_{Zb} \sigma_{ZW}}$$

is potentially competitive. The key is optimizing σ_{ZW} .

Analysis of σ_{ZW}

We have carried out a fast simulation of $e^+e^- \rightarrow e^+e^-h \rightarrow e^+e^-ZZ$ and backgrounds.

Events generated using the ILC-Whizard setup provided by Mikael Berggren. Beam profiles generated by GuineaPig, Whizard 1.95 matrix elements, Pythia showering and hadronization, SGV3.0 detector simulation.

Test results compare well for $Zh \rightarrow \nu\nu bb$ done in full detector simulation by H. Ono.

Focus on $ZZ \rightarrow 4j$. Fully reconstructible final state.

Features many kinematic features to discriminate against background.

Limiting factor is total initial size. $N = 473$. Cuts must be **efficient**.

Summary of Results

Merge down to 4 hadron jets.
Primary cuts

		Before Cuts		After Cuts	
		Zh	eeZ	$Z + X$	$ee + X$
① Opposite sign ee					
② $70(110) < M_{ee} < 110(150)$	Signal	158	315	85	183
③ $110 < M_{rec} < 250$	$eeqq$	931500		65	100
④ $M_{4j} < 150$	$eeWW(4j)$	8750		2	9
⑤ $W_{off} < 70, 55 < W_{on} < 100$	$eeWW(2j\nu)$	8390		1	9
	$eeZZ/\gamma\gamma$	21235		37	76
⑥ $p_W^{hrest} < 45$	$\delta\sigma(ee \rightarrow eeh, h \rightarrow WW \rightarrow 4j) = 8.8\%$				
⑦ $L(\theta_j, \dots, Y_{34}, Y_{45})$					

Adding in $Z \rightarrow \mu\mu$ conservatively

$$\delta\sigma_{ZW} = 7.2\%$$

Conclusions

- Using σ_{ZW} at 500 fb^{-1} , according to our simulations gives

$$\delta_{\Gamma} = 12\%$$

- Using full data from 250 and 500 yields

$$\delta_{\Gamma} = 6.5\%$$

- Combined with σ_{WW} data

$$\delta_{\Gamma} = 5.8\%$$

Conclusions

- ILC provides an ideal environment to assess Higgs width in model-independent way with high precision.
- 250 GeV run suggests multiple avenues to achieve best overall value.
- 500 GeV run can be dominated by a few well-studied modes. However, σ_{ZW} channels offer a competitive determination, and can be combined for best sensitivity.
- Further studies may bring more channels to prominence. Especially need improved $Zh \rightarrow ZZZ$ and LHC modes at 250.
- The future may hold surprises! All estimates based on SM-like branching fractions.